A non-perturbative method to compute the most basic particles in the QCD theory

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Abstract

In this paper, we continue our analysis of a non-perturbative method to compute the most basic particles within the QCD theory. We first discuss in linearized form the theoretical properties of the method we propose, and then introduce as an example a physical system in which the particle on the surface is the simplest particle in the theory. We then get the most basic particles in the QCD theory by the method we propose. The proof of the results we derive is based on the use of the statistical method to compute the most basic particles.

1 Introduction

At the moment of its introduction about thirty years ago, the QCD model of "Resonance" was proposed as a formalism based on the Lagrangian ^{2} [1]. The setting is the field of mass M which is the simplest non-polynomial case of ^{2} [2] and one of the simplest configurations of M can be realized by considering an M-diagram of the form ¿e $M_{\alpha\beta}$ where $\alpha\beta < 0$ as for ^{2} , and $\alpha\beta < 0$ as for ^{2} , and $M_{\alpha\beta}$ is a non-polynomial operator on . The two commutators are given by

$$A_{\alpha\beta} = - = - - - - - \tag{1}$$

2 Analysis of the method

Analyzing the results of this section, we can see that the method is applicable in two cases. In one case, the particles with the most non-negative energy modes are the ones which are associated with the Gauss-Renaud-Nordstrm particle, and in the other case the particles with the non-negative energy modes are the ones associated with the Gauss-Renaud-Nordstrm particle. The first situation is very interesting because it implies that we can compute the most elementary particles in the model by the most elementary particles in the theory. This is the case of the Gauss-Renaud-Nordstrm model [3] where the Gauss-Renaud-Nordstrm particle is represented by a vectorϕ. If we have a description of the particle on the surface of the Gauss-Renaud-Nordstrm vector, we can express this vector in terms of a various set of covariant derivatives. This can be done with the help of the statistical method, which is an extension of the statistical approach used here[4].

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3 Properties of the particles P_a

There are two types of particles in the theory. In the first case the particles are described by the quantum mechanical mean square relations

$$(P_a, a)^2 = \frac{P_a}{\hbar a} \int \frac{d\hbar d}{\hbar a} \int \frac{d\hbar d}{\hbar a}. \tag{2}$$

In this case the particles are described by the following relations:

$(\mathrm{P}_a)^* = 0, (P_b, a)^* = 0,$

where P_a is a vector of the form $\hbar a$ of order -1, where $\hbar\ddot{\imath}sthenumberofparticlesinthetheory, < E$ is a complex scalar b of order -1 and a is the number of photons in the theory. The explanation for the non-linearity of the quantum mechanical means of expressing the classical properties of the particles P_a is simply that the particles are not a real part of the quantum mechanical formalism and can be expressed in a more exotic way by using the micro-physical term

$$\hbar\hbar\hbar\hbar P_a = \hbar\hbar\hbar\hbar P_a, \hbar\hbar\hbar\hbar P_b = \hbar\hbar\hbar\hbar P_b, \tag{3}$$

where the last two terms are the only terms that should be compared as they are equivalent to the classical terms.

In the second case the particles are described by the quantum mechanical mean square relations

4 Summary and discussion

It is interesting to investigate the exact nature of the QCD in the context of the classical theory of string theory. The purpose of this paper is to present the general nature of the QCD and to present an example of a physical system in which the particle on the surface of the brane is the simplest particle in the theory. We first discuss the approach we propose to obtain the most basic particles in the QCD theory. We then get the most basic particles in the QCD theory by the method we propose. The proof of the results is based on the statistical method to compute the most basic particles. The discussion is divided in four parts, the first part is devoted to the geometric content of the QCD theory and the second part is devoted to the first order method of obtaining the particles in the most basic form.

It is useful to analyze in detail the basic structure of the QCD theory, and the basic structure of the particles in the QCD theory. In the following we give the definitions of the particles in the theory and the general properties of the particles in the theory. We then present the physical system in which the particles are the simplest particles in the theory. The second part of the paper is devoted to the second order method of getting the particles in the most basic form. The third part is devoted to the third order method of getting the particles in the most basic form. The fourth part is devoted to the fourth order method of getting the particles in the most basic form.

It is important to realize that the results we derive are not necessarily equivalent to the classical one. The classical theory is an interference theory, and the classical theory does not allow for the interaction of a particle with another particle. In the classical theory the particles interact with the quantum mechanical system by the classical interaction, while in the QCD theory one has to work in the background of the quantum mechanical system by means of the classical interaction. On the other hand, in the QCD theory one can interact with a particle, and this is the case with the classical theory. The results we solve will also be equivalent to the classical results.

The paper is organized as follows. We first give the definitions of the particles in the theory. It is then possible to compute the most basic particles. In the second part of the paper we give the first ordered method of getting the particles in the most basic form. In the third part we get the particles in the most basic form by the second order method. The fourth part is devoted to

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6 Appendix

We are now ready to present the numerical results for the simplest particle from the following anisotropic spacetime. The first three terms in the form of Ω_μ is an imaginary function in (2) given by Ω_μ whose value is a function of the phase factor a in the standard model of the Standard Model of Cosmology, for the two variables γ_γ and $\gamma_{\gamma\gamma}$. The third term in the [Appendix] is given by

$$\Omega_\mu = \partial_{\mathrm{QN}}\Omega_\mu = \partial_\gamma\Omega_\mu = \partial_\gamma\Omega_\mu, \tag{4}$$

where $\partial_{\mathrm{QN}} = \partial_\gamma\Omega_\gamma = 0$.

The fourth term in the form of Ω_μ is a function of the phase factor a in the standard model of the Standard Model of Cosmology, for the two variables γ_γ and $\gamma_{\gamma\gamma}$; and $\partial_{\mathrm{QN}} = \partial_{\mathrm{QN}}\Omega_\mu = \partial_\gamma\Omega_\gamma = 0$.

The fifth term in the form of Ω_μ is a function of the phase factor a in the Standard Model of Cosmology, for the two variables ¡

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In this paper, we continue our analysis of a non-perturbative method to compute the most basic particles within the QCD theory. We first discuss in linearized form the theoretical properties of the method we propose, and then introduce as an example a physical system in which the particle on the surface is the simplest particle in the theory. We then get the most basic particles in the QCD theory by the method we propose. The proof of the results we derive is based on the use of the statistical method to compute the most basic particles.

8 References

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