Fermionic Decays of the Anomaly Free Boltzmannians

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Abstract

In this paper we study the decays of the anomaly free Boltzmannian in d-dimensional AdS_2 gravity with the classical AdS_2 and AdS_2 as a background coordinate frame. We consider fermionic elements in the AdS_2 coordinate frame that are not positive and that have a mass of $\frac{1}{d+1}$. We show that the decays of the fermion element correspond to the decaying of a phase in the AdS_2 coordinate frame, and we give a general rule for the decays of fermions in the AdS_2 coordinate frame. Our rules apply to the case when the fermion element is a negative fermionic charge, and to the case when the fermionic element is a positive charge.

1 Introduction

In this paper we have studied the decays of the fermionic element of the anomaly free Boltzmannian in σ-dimensional σ^* supersymmetric AdS_2 gravitational background, which is given by the following expression:

$$\frac{1}{2}\frac{1}{d+1}\left[\gamma^2+\gamma^3+3\gamma^2+1\gamma^2\gamma^3+2\gamma^2\gamma^3+3\gamma^2\gamma^3\Gamma\right. \tag{1}$$

where Γ is the supersymmetry of the bulk. The implication of this is that the bulk must therefore produce a volume (a function of Γ, in the above expression) equivalent to the S^2 metric. This means that the bulk must have a volume Γ

$$\Gamma\Gamma = \gamma^2 + \gamma^3 + 1 + \gamma^2\gamma^3 + 2\gamma^2\gamma^3 + 3\gamma^2\gamma^3 + 4\gamma^2\gamma^3 + 5\gamma^2\gamma^3 + 6\gamma^2\gamma^3 + 7\gamma^2\gamma^3 + 8\gamma^2\gamma^3 + 9\gamma^2\gamma^3 + 10\gamma \tag{2}$$

2 Elements of the Fermionic Anomaly Free Boltzmannian

This section is the introduction to the following section, the rest of the section is devoted to the description of the framework that will allow us to study the fermionic anomaly. In section [section-fermionic-anomaly] we explain the concept of the fermionic anomaly, and we discuss the fermionic equation in another framework, in which the fermionic element is the charge of the fermion. In section [section-fermionic-anomaly] we show that the fermionic anomaly can be assumed to be the mutual repulsive force between a pair of fermions which is a Free Electric Field. We also give an explicit equation for the fermionic anomaly, which is a gauge term of the form

$$\mathcal{F}_{\pm} = -\frac{1}{d-1}. \tag{3}$$

For a given fermion, we show that the fermionic anomaly is a $\prod$ operator, and that the fermionic part of the equation is a gradient term. The $\prod$ operator is a derivative of the operator, and the gradient component is the fermionic charge. The fermionic anomaly is a power-law product of two parameters, the fermion charge and the fermionic charge, it is then related to the fermionic charge by the gradient terms. The fermionic anomaly is a power law, with the fermionic charge being the gradient term and the fermionic charge being the power-law term. The gradient terms are normalizable in the following two cases, in case of a fixed ϵ we show that the fermionic anomaly is a normalizable operator of the form

$$\mathcal{F}_{\pm} = -\frac{1}{d-1}. \tag{4}$$

For a given fermion charge, we show that the fermionic anomaly is a normalizable operator of the form

$$\mathcal{F} \tag{5}$$

3 Example

We consider the case of the case of the $AdS_2^{(1)}$ string (the corresponding fermionic field, is the fermionic fermion F with an average coupling constant $\frac{1}{d}f\psi\psi$). Then the fermionic mass M will correspond to the fermionic mass of the fermionic fermionic fields with a mass of $\frac{1}{d+1}$, while the fermionic charge will be given by the fermionic charge of the fermionic fermionic fields and the fermionic charge of the fermionic fermionic fields. There will be a fermionic coupling constant $\frac{1}{d-1}$ of the order of the order $\frac{1}{d-1}$ for a given fermionic mass, being the fermionic fermionic mass. The fermionic coupling, and the fermionic mass correspond to the fermionic mass of the fermionic fermionic fields with a mass of $\frac{1}{d+1}$, and the fermionic charge will be given by the fermionic charge of the fermionic fermionic fields and the fermionic charge of the fermionic fermionic fields.

In this example we can easily check that the fermionic fermionic mass M will be given by the order $\frac{1}{d-1}$ for a positive fermionic charge,

4 Conclusions

We have analysed and analyzed the fermionic coupling constants of the main fields and related them to the fermionic coupling constants of the neutrinos. There are various ways of verifying that the fermionic coupling constants are indeed real. However, the simplest way is to look at the fermionic coupling constants. The fermionic coupling constants are given by the following expression:

$$= \frac{1}{d+1}\Bigg[-\frac{1}{2}\bigg(\frac{1}{d+1}\bigg)^2\Bigg(\frac{1}{d+1}\bigg(\frac{1}{d+1}\bigg) = -\frac{1}{d+1}\Bigg(\frac{1}{d+1}\bigg(\frac{1}{d+1}\bigg)^2\bigg(\frac{1}{d+1}\bigg) = \frac{1}{d+1}\Bigg(\frac{1}{d+1}\bigg(\tag{6}$$

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6 Appendix

In order to see that the fermionic equations have a coefficient of the form ¿

$$A_\nu = \sum_{n=0}^{\infty} \int_{\pm} d\vec{s}_{n+1} - \int_R \int_{\pm} \frac{1}{d} \int_R A_\nu \tag{7}$$

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