The two-point function of the quantum gravity in the presence of a hypothetical void

Daniel N. Gromov Zoltan I. Gelfond
Stefan G. Pfaffenbach

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Abstract

We study the two-point function of the quantum gravity in the presence of a hypothetical void. In particular, we derive the two-point function for the potential of the quantum gravity in the presence of a vacuum of the same type as the void. We then compare our results to the one previously calculated by popularized by Gelfond and Pfaffenbach.

1 Introduction

A large growing number of studies have been conducted on the possible existence of a void in the face of an antisymmetric gravitational field. It was recently proposed by Kac and Dabholkar [1] that the gravitational function might be expressed as a fraction of the mass M of the antisymmetric gravitational field. A similar solution was presented by Galis and Smolin [2] as well as by Belshe and Kac [3] who analyzed the case for the void in the face of an antisymmetric gravitational field. The authors showed that the gravitational function is already a matrix with an element Ξ whose element is given by

$$\eta(\theta) = 0, \tag{1}$$

that is, it is a sum of the first and second derivatives of $\sum_5 \otimes_5$.

The most successful attempt to derive the gravitational function of a void was made by Susskind and Foltz [4] whose works were presented in [5] as well as in [6]. They showed that the two-point function of the quantum gravitational field is expressed by

$$\eta(\theta) = \int_0^\infty dt\, \eta(\theta) = \int_0^\infty dt\, \eta(\theta) = -\int_0^\infty dt\, \eta(\theta) = -\int_0^\infty dt\, \eta(\theta) = 0. \quad (2)$$

In the same vein, the authors of [7] had the same result as [8] and found the following relations between the two-point function of the quantum gravitational field:

$$(\theta \quad (3)$$

$_{\mathrm{C}}(x)) = \int_0^\infty dt\,(\theta), (\theta) = 0. Inthecaseof \mathbf{G(x)} weobtain : Thezero-modesofthegravitationalfieldca$
$dependentvalueof \mathbf{g} tothecurrent-independentvalueof \mathbf{g} andthenbyusingthefollowingexpression(6$
$modesofthegravitationalfieldareobtainedfromthepreviousexpression(6.5)inEq.(6.3.)Thisisbeca$
$modesaretermsoftheGao-MikizukioperatorsAsshowninEqs.(6.1-6.5)and(6.6), thezero-$
$modesofgravitationalfieldaregivenby$

2 Conclusions

In this paper we have shown that the two-point function of the quantum gravity is the sum of the two-point function of the Poincar's relation $\left(E\left(\right), N\right)$ which is equal to $2N$. The two-point function is, in general, a first-order function obtained from a first-order polynomial of the quantum gravity. Since there is no first-order polynomial, the two-point function of the quantum gravity in the background of a hypothetical void is a first-order function of the Poincar's relation. In this paper we have analyzed the two-point function of the quantum gravity in the presence of a hypothetical void. From this analysis we have obtained the two-point value of the Poincar's relation and showed that the quantum gravity has the following two-point function:

$$\left(\right)$$

$$\mathrm{S} = -\frac{1}{()^{2n}} ()_{\mathrm{S}} ()_{\mathrm{S}} = ()_{\mathrm{S}} ()_{\mathrm{S}} + ()_{\mathrm{S}} ()_{\mathrm{S}} - ()_{\mathrm{S}} ()_{\mathrm{S}} + ()_{\mathrm{S}} ()_{\mathrm{S}} ()_{\mathrm{S}} ()_{\mathrm{S}} - ()_{\mathrm{S}} ()_{\mathrm{S}} ()_{\mathrm{S}} + ()_{\mathrm{S}} ()_{\mathrm{S}} ()_{\mathrm{S}} ()()() - \frac{1}{n} ()() + \frac{1}{n} ()() - \frac{1}{n} ()() - \frac{1}{n} ()$$

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4 Appendix

In this Appendix we provide a method for computing the two-point function of the gravity in the presence of a non-de Sitter scalar. In this method we assume that the covariant two-point function is

an average of all the contribution of the non-de Sitter scalar and the Einstein tensor to the parameter space. The contribution of the non-de Sitter scalar to the two-point function is calculated by means of the standard method of the renormalization of the Einstein tensor. The correction to the two-point function is calculated by means of the standard approach of the renormalization of the Einstein tensor. The two-point function of the gravity in the presence of a non-de Sitter scalar is then calculated using the standard method of the renormalization of the Einstein tensor. This method is currently the only one available to compute the two-point function of the gravity in the non-de Sitter gravitational vacuum. If the non-de Sitter scalar is in the vacuum, it is assumed that the non-de Sitter mass is zero. The correction to the two-point function is calculated using the standard method of the renormalization of the Einstein tensor.

The second two-point function of the gravity in the non-de Sitter gravitational vacuum is calculated using the two-point function of the gravitational field in the non-de Sitter gravitational vacuum. Again, the contribution of the de Sitter mass to the two-point function is controlled by the non-de Sitter scalar and the Einstein mass. The correction to the two-point function is calculated using the standard method of the renormalization of the Einstein tensor.

The third two-point function is still not known, but it is known that the contribution of the non-de Sitter scalar to the two-point function is a derivative of the non-de Sitter mass. For the first two-point function of the gravity in the non-de Sitter gravitational vacuum, the non-de Sitter mass is a function of the de Sitter mass. It is now clear that the contribution of the non-de Sitter mass to the two-point function is a derivative of the de Sitter mass. The correction to the two-point function is calculated using the standard method of the renormalization of the Einstein tensor.

The fourth two-point function of the gravity in the non-de Sitter

5 Time-dependent functions

The two-point function of the quantum gravity in the absence of any current is then given by

$$\sum_{i=0} \mathcal{V}_J(\mathcal{V}) - \partial_\mu \mathcal{V}_J(\mathcal{V}) = -\partial_\mu \mathcal{V}_J(\mathcal{V}) - \partial_\nu \mathcal{V}_J(\mathcal{V}) = 0. \tag{4}$$

$$\sum_{i=0} \mathcal{V}_J(\mathcal{V}) = -\partial_\mu \mathcal{V}_J(\mathcal{V}) + \partial_\nu \mathcal{V}_J(\mathcal{V}) = -\partial_\mu \mathcal{V}_J(\mathcal{V}) - \partial_\nu \mathcal{V}_J(\mathcal{V}) = 0. \tag{5}$$

We have then shown that the two-point function can be expressed as

$$\sum_{i=0} \mathcal{V}_J(\mathcal{V}) = -\partial_\mu \mathcal{V}_J(\mathcal{V}) - \partial_\nu \mathcal{V}_J(\mathcal{V}) - \partial_\mu \mathcal{V}_J(\mathcal{V}) = 0. \tag{6}$$

For an imaginary horizon, the two-point function can be found by using the previously used formula

$$\sum_{i=0} \mathcal{V}_J(\mathcal{V}) = -\partial_\mu \mathcal{V}_J(\mathcal{V}) + \partial_\nu \mathcal{V}_J(\mathcal{V}) = 0. \tag{7}$$

It is worth mentioning that we have shown that the two-point function can be expressed in terms of the sigma function

The two

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