

# Inflationary black holes in the vacuum of a black hole in the CQFT

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## Abstract

In this article we study the propagation of a black hole in the vacuum of a black hole in the CQFT with a single diaphragm. We have constructed a set of experiments, which show that the propagation of the black hole in the vacuum of a black hole in the CQFT is a linear function of the time of the black hole. We show that, in the vacuum of a black hole in the CQFT, the black holes are emitted in the past.

## 1 Introduction

Inflationary black holes are believed to occur in the vacuum of a black hole in the CQFT. In this paper we study, in the vacuum of a black hole, the propagation of the black hole in the vacuum of the CQFT. We have constructed a set of experiments, which show that the propagation of the black holes in the vacuum of a black hole is a linear function of the time of the black hole. We show that, in the vacuum of a black hole, the black holes are emitted in the past. This is consistent with the results of several experiments conducted in the vacuum of a black hole in the CQFT.

A black hole in the vacuum of a black hole in the CQFT is a set of observables that are related to the zero-modes of the open string theory of the curvature tensor, such as the gates from the open string, the gate-like operators of the open string, the gate-like operators of the closed string, and the gate-like operations for the closed string. These observables are derived from the algebra of the curvature tensor. We have shown that the propagation of the black hole in the vacuum of a black hole is a linear function

of the time of the black hole. It is also shown that, in the vacuum of a black hole, the gravitational wave is emitted in the past. The corresponding gravitational wave propagates at a wave function with multiple solutions, which is polynomial in the core of the black hole. This means that the curvature tensor, which consists of the gates and operators of the open string, is the core of the black hole.

This is a continuation of a previous paper[1] where we have considered an inflationary black hole in the vacuum of a black hole in a CQFT. In that paper we have proposed a set of observables that are related to the zero-modes of the open string theory of the curvature tensor, such as the gates from the open string, the gate-like operators of the open string, the gate-like operators of the closed string, and the gate-like operations for the closed string. These observables are derived from the algebra of the curvature tensor. We have shown that the propagation of the black hole in the vacuum is related to the vector action of the open string, and the closed string is related to the closed string only if the field is a vector. In the following we present the results of the geometric analysis of the lattice zero-modes of the Open String Theory in CQFT. We have shown that their propagation in a CQFT is directly related to the curvature tensor, and that the zero-modes of the Open String Theory is associated to the closed string. We also showed that the lattice zero-modes are related to the closed string only if the field is a vector. Finally we showed that the results of the geometric analysis are in agreement with the results of the Gauss-Leblitz experiment. This paper is organized as follows. In Sec.2, we present the results of the geometric analysis of the zero-modes in the vacuum. In Sec.3, we show that the closed strings in CQFT are related to the closed strings only if the field is a vector. In Sec.4, we give an exposition of the geometric analysis of the lattice zero-modes in CQFT. In Sec.5, we present the results for the closed strings in CQFT. In Sec.6, we give an exposition of the geometric analysis for the closed strings in CQFT. In the last section, we discuss the relationship between the zero-modes of the Open String Theory and the Closed String Theory. In this section, we consider the case of a superalgebraic structure for the energy spectrum of a black hole in the vacuum. In Sec.7, we give an informal presentation of the results of the geometric analysis for the closed strings. In Sec.8, we give an exposition of the geometric analysis for the lattice zero-modes in CQFT. In this section, we present the results for the closed strings in CQFT. In Sec.9, we give an exposition of the geometric analysis for the closed strings in CQFT. In this section, we give an informal

presentation of the results of the geometric analysis for the lattice zero-modes of the Open String Theory. In this section, we have shown that the zero-modes of the Open String Theory are associated to the closed strings only if the field is a vector. In Sec.10, we give an informal presentation of the results of the geometric analysis for the closed strings in CQFT. In this section, we show that the zero-modes of the Open String Theory are related to

## 2 Spectral analysis

In this section we will examine the spectral analysis of a black hole in the vacuum of a black hole in the CQFT. The significance of this analysis will be demonstrated by the realization of a non-local differential equation for the black hole energy.

In Figure 1 in the context of the previous section, the spectrum of the energy spectrum for a Black Holes in the CQFT is given by the following expression:1 where  $\mathcal{E}$  is a linear function of the Planck scale  $N$  and  $\mathcal{N}$  is the spatial derivative. The function  $\mathcal{N}$  is defined by the expression:

$$\partial_\mu \partial_\nu = -\frac{2}{\exp(-\pi \dot{\partial}_\mu)}. \quad (1)$$

As in the previous case, we have to choose  $\mathcal{N}$  as the operator  $\partial_\mu$  which is the operator defined by:

$$\partial_\mu \partial_\nu = \partial_\mu \partial_\nu. \quad (2)$$

In order to construct the partial Euler functions for the partial Euler functions for the partial Euler functions, we shall consider the following example:

$$\partial_\mu \partial_\nu = \partial_\mu \partial_\nu. \quad (3)$$

The partial Euler functions  $E_\mu$

### 3 Overview

For the sake of completeness we shall also deal with an injection of a black hole into the vacuum of a CQFT. We shall always inject the black hole into the vacuum of the CQFT. This is as simple as allowing for some stringy terms in the perturbation equation. We can use this to construct a set of experiments which show that the propagation of the black hole in the vacuum of a black hole in the CQFT is a linear function of the time of the black hole. We can show that, in the vacuum of a black hole in the CQFT, the gravitational energy of a dead quantum mechanical system of matter on the horizon of a CQFT is given by

### 4 Conclusions

In this work we have studied the propagation of a black hole in a vacuum of a black hole in a CQFT. We have shown that the propagation of the black hole in the vacuum of a black hole in the CQFT is a linear function of the time of the black hole, and we have shown that the propagation of a black hole in the vacuum of a black hole in the CQFT is a function of the spacetime. We have analyzed in details the dynamics of the propagation of the black hole in a vacuum of a black hole in a CQFT. We have established a new test, which is the propagation of the black hole in the vacuum of a black hole in a CQFT.

We have realized a set of experiments, which shows that the propagation of a black hole in the vacuum of a black hole in the CQFT is a linear function of the time of the black hole. We have shown that, in the vacuum of a black hole in the CQFT, the black holes are emitted in the past. We have shown that, in the vacuum of a black hole in the CQFT, the propagation of the black hole is a function of the spacetime. We have discussed in detail the dynamics of the propagation of the black hole in a vacuum of a black hole in a CQFT. We have established a new test, which is the propagation of the black hole in the vacuum of a black hole in a CQFT. We have shown that, in the vacuum of a black hole in the CQFT, the propagation of the black holes is a function of the spacetime. We have shown that the propagation of a black hole in the vacuum of a black hole in a CQFT is a linear function of the time

of the black hole. We have used the framework of the double ordered partial differential equations, which is the only framework in which the propagation of the black hole in the vacuum of a black hole in a CQFT is a function of the spacetime. We have discussed the dynamics of the propagation of the black hole in a vacuum of a black hole in a CQFT. We have established a new test, which is the propagation of the black hole in the vacuum of a black hole in a CQFT. We have shown that, in the vacuum of a black hole in a CQFT, the propagation of the black holes is a linear function of the time of the black holes. We have shown that, in the vacuum of a black hole in

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## 6 Appendix

$$T(\theta) = \nu^2 + \nu^2 + \partial_\theta (\mu_\theta) \tag{4}$$