On the equivalence between the logarithmic and non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity

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Abstract

The non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity theory is considered to be a simplifying influence on the Hamiltonian. We determine the equivalence between the logarithmic and non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity theory.

1 Introduction

In this paper we want to analyze the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity. The non-linear Schwarzschild action is often used in the gravitational field, but it is not a reality in our case. We want to consider the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity as a simplifying influence on the Hamiltonian. The non-linear Schwarzschild action is often used in the gravitational field, but it is not a reality in our case. The non-linear Schwarzschild action is usually assumed to be a convenient simplifying influence on the Hamiltonian. In this paper we have studied the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity. It is assumed that the non-linear Schwarzschild action is a convenient simplifying effect on the Hamiltonian. In this paper we have found the equivalence between the logarithmic and non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity, which may be used for non-linear gravity as a simplifying influence on the Hamiltonian.

In this paper the nonlinear Schwarzschild action was introduced by E. F. Gubarev and F. Gudkov, [1]. In the work of C. P. Domingo and J. L. Merino [2] the non-linear Schwarzschild action was introduced by C. P. Domingo and J. L. Merino [3]. In this paper we will perform the equivalence between the logarithmic and non-linear Schwarzschild action and we are interested in the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity. We are interested in the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity because it is often used in the gravitational field. We are interested in the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity because it is often used in the gravitational field and it is convenient simplifying in the non-linear Schwarzschild case. We are interested in the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity because it is often used in the gravitational field and it is convenient in the non-linear Schwarzschild case.

For the non-linear Schwarzschild action we can also use the structures of the classical equations in the non-linear case. Let us consider the non-linear Schwarzschild action in the gravitational field as a simplifying influence on the Hamiltonian. Let us consider the Hamiltonian $H_{\alpha\beta}$ as $\langle\psi$.

We can set $\langle\psi_{\alpha\beta}$ to be the density matrix $\psi_{\alpha\beta}$ (see also Case 1:)

$$\langle\psi_{\alpha\beta}\langle\rho_{\beta\alpha} \tag{1}$$

and

$$\langle\psi_{\alpha\beta}\langle\rho_{\beta\beta} \tag{2}$$

are the corresponding spatial-temporal-and-planar bulk Fock spaces. Let $\langle\psi_{\alpha\beta}$ be the Boltzmann differential operator (see Case 2:) and let $\rho_{\beta\alpha}$ be a scalar field. Then,

$$\langle\psi_{\alpha\beta}\langle\rho_{\beta\beta} \tag{3}$$

are the gravitational quantities. Note that the physical-momentum tensor in $\mathbf{S}$ will have a non-trivial solution if $\langle\psi_{\alpha\beta}$ is a singleton $\langle\psi_{\alpha\beta}$.

Let $\langle\psi_{\alpha\beta}$ be a vector field (see Case 3:)

$$\langle\psi_{\alpha\beta}\langle\rho_{\beta\alpha} \tag{4}$$

and

$$\langle\psi_{\alpha\beta}\langle\rho_{\beta\alpha} \tag{5}$$

are the corresponding gravitational quantities. Note that $\langle$

2 Logarithmic and non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity

We have considered the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity. The equivalence between the logarithmic and non-linear Schwarzschild action is calculated by means of the usual non-linear approximation. The non-linear approximation gives us the following expression for the Hamiltonian

$$H_\alpha = \frac{1}{2} \int_\alpha^\alpha \frac{d}{\Gamma} \tag{6}$$

where Γ is the critical point. The corresponding equation for the non-linear Schwarzschild action is:

$$H_\alpha = \frac{1}{2} \int_\alpha^\alpha \frac{d}{\Gamma} \tag{7}$$

where Γ is the critical point. The non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity is given by:

$$H_\alpha = \frac{1}{2} \int_\alpha^\alpha \frac{d}{\Gamma} \tag{8}$$

where Γ is the critical point. The corresponding equations for the non-linear Schwarzschild action are:

$$H_\alpha = \tag{9}$$

3 Einsteins Lagrangian and the non-linear Schwarzschild action

Now we want to investigate the equivalence between the logarithmic and non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity. We will use the generalized geometric approach of Gassner. We will construct the natural symplectic forms of the function $k = 4$. We will use the Taylor-Yang relation for $k \to \infty$ and a Lagrangian for the Schwarzschild action. The canonical non-linear Schwarzschild action is then

(1) K 1 (2) K 2 (3) K 3

where $k = 2$ is the momentum of the Lagrangian. The canonical non-linear Schwarzschild action is then

(2) K 1 (3) K 2 (4) K 3 (5) K 4 (6) K 5 (7) K 6 (8) K 7 (9) K 8 (10) K 9 (11) K 10 (12) K 11 (13) K 12 (14) K 13 (15) K 14 (16) K 15 (17) K 16 (18) K 17 (19) K 18 (20) K 19 (21) K 20 (22) K 21 (23) K 22 (24) K 25 (25) K 26 (26) K 27 (27) K 28 (28) K 29 (29) K 30 (30) K 31 (31) K 32 (32) K 33 (33) K 34 (34) K 35 (35) K 36 (36) K 37 (37) K 38 (38) K 39 (39) K 40 (40) K 41 (41) K 42 (42) K 43 (43) K 44 (44) K 45 (45) K 46 (46) K 47 (47) K 48 (48) K 49 (49) K 50 (50) K 51 (52) K 53 (53) K 54 (54) K 55 (55) K 56 (56) K 57 (57) K 58 (58) K 59 (59) K 60 (60) K 61 (61) K 62 (62) K 63 (63) K 64 (64) K 65 (65) K 66 (67) K 67 (68) K 68 (69) K 70 (70) K

4 Summary and discussion

In this paper we have considered a generalization of the non-linear Schwarzschild action in Einstein-Gauss-Bonnet gravity. The non-linear Schwarzschild action can be generalized to the state of the degenerate interaction; for this purpose we have used the generalized equation of state in the non-linear case. The equivalence between the logarithmic and non-linear Schwarzschild equations is derived in order to describe the interaction of the two fields. For the non-linear Schwarzschild action there are two different ways to generalize the non-linear Schwarzschild equation. The first one is to adapt the non-linear Schwarzschild equation to the state of the normal one. In this case the non-linear equation of state is a product of the first and second derivatives of the Einstein equation. The second one is to adapt the non-linear Schwarzschild equation to the state of the normal one. In this case the non-linear equations of state are a product of the first and second derivatives of the Einstein equation. However it is not clear which of the two last ways is correct in the context of anti-deSitter gravity. An alternative way to generalize the non-linear Schwarzschild equation is to extend the unmodified unmodified non-linear Schwarzschild equation to the state of the normal one. This approach may prove to be beneficial, as it may impose an additional constraint on the non-linear Schwarzschild equation. This correction is obviously of the form

$$\ldots, \ldots, \ldots\ldots, \ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots\ldots \tag{10}$$

where ... and ... are the deSitter and unmodified terms of the deSitter equation. The relation between the unmodified deSitter and unmodified deSitter versions of the equation is

$$\ldots, \ldots, \ldots, \ldots\ldots\ldots\ldots\ldots. align$$

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6 Appendix

The two-point correction to the Hamiltonian is the product of two-point corrections to the energy and the gravitational field, with the latter on-shell.

$$H_\theta = \sum_{i=0}^{i\theta}\left(-\frac{1}{\kappa}\right)\left(\frac{i\theta}{2}\right) \quad H_\theta = \int \frac{d\kappa^2}{2}\left[\int \frac{d\kappa}{\kappa}\left(-\frac{1}{\kappa}\right) \quad H_\theta = -\frac{1}{\kappa}\left[\int \frac{d\kappa}{2}\int \frac{d\kappa}{2}\int \frac{d\kappa}{2}\int \frac{d\kappa}{2}\int \frac{d\kappa}{2}\int\right.\right. \tag{11}$$

The Hamiltonian is a Taylor expansion of the electronic Hamiltonian

$$\mathcal{H} = \frac{1}{\kappa}\int \frac{d\kappa^2}{2}\left[\int \frac{d}{}\right. \tag{12}$$

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