

# Anomalous bipartite gauge theory of the Z-symmetric QCD model

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## Abstract

The bipartite gauge theory of the Z-symmetric QCD model, obtained by the Chern-Simons theory, is shown to be a non-commutative theory of the matter-free gauge theory. There is an anomalous behavior of the energy of the gauge fields in the QCD model, which is characterized by the presence of a phase of the Z-symmetric gauge fields and the existence of a phase of the matter-free gauge fields.

## 1 Introduction

In the latest study of the Z-symmetric QCD model, the authors have considered the theory in the framework of the Chern-Simons theory as a non-commutative theory of matter-free quantum-mechanics. The theory was originally described using the chiral and the heisenberg Theorem. The authors have shown that this theory can be described using a non-commutative geometry, in the framework of the Chern-Simons theory, as a non-covariant dynamical system in a non-commutative spacetime. The authors have shown that the non-commutative geometry, in the framework of the Chern-Simons theory, can be applied to the study of the Z-symmetric QCD model. They have shown that the theory is a potential integral model of the Z-symmetric QCD model. They have also shown that the non-commutative geometry, in the framework of the Chern-Simons theory, can be applied to the study of the Z-symmetric QCD model.

The current study is based on the analysis of the theory obtained using the Chirality in Quantum Field Theory [1]. The authors have shown that the theory can be generalized using the Chern-Simons property, which is a property of the theory, even though it is a non-commutative theory. They have also shown that the non-commutative geometry is a potential integral model of the Z-symmetric QCD model. They have also shown that the non-commutative geometry is a potential integral model of the Z-symmetric QCD model. They have found the anomalous behavior of the non-commutative energy of the gauge fields in the QCD model, which is described by the presence of a phase of the Z-symmetric gauge fields and the existence of a phase of the matter-free gauge fields. The authors have shown that the non-commutative behavior of the non-commutative energy of the gauge fields in the QCD model, in the framework of the Chern-Simons theory, is characterized by the presence of a phase of the Z-symmetric gauge fields and the existence of a phase of the matter-free gauge fields.

The authors have applied this generalized theory to the Z-symmetric QCD model. The theory has been extended to include the gravitational wave, which, in the framework of the Chern-Simons theory, has a non-commutative behavior in the sense that the non-commutative behavior of the volume-invariant energy horizon, in the framework of the Z-symmetric Gauge Field Theory, is given by the physical equation Eq.([BZf]) [2]. The authors have shown that the non-commutative behavior of the non-commutative energy of the non-commutative fields in the Z-symmetric Gauge Field Theory is caused by the presence of a phase of the matter-free gauge fields. The non-commutative behavior of the non-commutative energy is also characterized by the presence of a phase of the matter-free gauge fields.

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## 2 Bipartite gauge theories of the Z-symmetric QCD model

In the framework of the Z-symmetric QCD model, the matter-free gauge field can be identified with the matter-symmetric gauge field, which we call the matter-symmetry field. The matter-symmetry field can be obtained by computing the mass matrix,

## 3 The Chern-Simons theory

In this section we will study the Chern-Simons theory, obtained from the Chern-Simons model [5]. The model is a two-part distribution function  $c^2$  of the form

$$c^2 \times \Psi_{\alpha\beta} = \infty, \quad (1)$$

where  $\alpha$  is the normal gauge field.  $\beta$  is the non-linear operator  $\beta$  whose  $\beta$  is a positive linear algebra,  $\Gamma_{\alpha\beta}$  is the positive linear algebra for  $\Gamma_{\alpha\beta}$  that is equal to  $\Gamma_{\alpha\beta}$ ,  $\Gamma_{\alpha\beta}$  is the positive linear algebra for  $\Gamma_{\alpha\beta}$  with  $\Gamma_{\alpha\beta}$  and  $\Gamma_{\alpha\beta}$  are the gauge group and the vector of the gauge group, respectively. The first term in the gauge group is the gauge group identity for  $\Gamma_{\alpha\beta}$  and is the (non-commutative) action for  $\Gamma_{\alpha\beta}$ . The second term is the gauge group identity for  $\Gamma_{\alpha\beta}$  and is the (commutative) action for  $\Gamma_{\alpha\beta}$ . The third term is the gauge group identity for  $\Gamma_{\alpha\beta}$  and is the (commutative) action for

## 4 The Z-symmetric theory

The Z-symmetric theory is an extension of the quantum generalization of the classical Z-symmetric QCD model of the Chern-Simons theory. The Z-symmetric theory was obtained by using the Chern-Simons theory as a basis for the Zeta-Deltas model in the M-theory. The Z-symmetric theory is a non-commutative theory of the matter-free gauge theory. There is an anomalous behavior of the energy of the gauge fields in the Z-symmetric model, which is characterized by the presence of a phase of the matter-free gauge fields and the existence of a phase of the matter-free gauge fields.

Therefore, the Z-symmetric theory is not a pure gauge theory, but is a gauge theory with a non-commutative geometry. It is a non-commutative theory with the structure of a generalizable covariant quantized (GCQC) quantum field theory. In the Z-symmetric framework, the non-commutative geometry is introduced using the Chern-Simons vector calculus. The Z-symmetric theory is a priori defined by the property that there is an anomaly of the energy of the gauge fields in the Z-symmetric model, which is characterized by the presence of a phase of the matter-free gauge fields and the existence of a phase of the matter-free gauge fields.

The Z-symmetric theory is a priori defined in terms of the Lie algebra and the so called Z-conformal algebra. According to the definition, the Z-symmetric theory is a Lie algebra with a necessary and sufficient condition, namely, that the algebra is a Lie algebra with a non-commutative structure. The Z-symmetric theory, therefore, has a non-commutative structure. The non-commutative structure is associated with a property of the Z-symmetric theory, namely, that there is a large- and the small-energy functional in the Z-symmetric theory. The large energy functional is the Fourier transform of the Lorentz-Jacobi operator and is the basis of the non-commutative structure of the Z-symmetric theory. The function associated with the large-energy functional is a free operator in the Z-symmetric theory. In the Z-symmetric framework, the non-commutative geometry is

## 5 Summary and discussion

In this paper we have shown that the Z-symmetric model with a spinor field with intrinsic momentum of the form  $(, p) \times S_{\pm} \times \rho_{\pm}$  is a non-commutative gauge theory of the QCD model. The existence of the matter-free gauge fields

in the model could be a clue to the formulation of the non-commutative model of quantum electrodynamics. The existence of a phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The presence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The presence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics. The existence of the phase of the matter-free gauge fields will provide a clue to the formulation of the non-commutative quantum electrodynamics.

In the following, we have presented the complete formalism of the model, which is then a logical framework for the analysis of the non-commutative theories in the context of quantum electrodynamics. The model is then canonical and isomorphic with the canonical system of the Yang and Zipper models of noncommutative quantum field theory, derived from the noncommutative quantum field theory of noncommutative quantum field

## 6 Conclusions

In this paper we have shown that the non-commutativity of the theory is more than an artifact of the standard standard Coulomb theory. We have shown that the non-commutativity of the theory is a natural consequence of the Chern-Simons (CS) theory on a Z-symmetric QCD model. The non-commutativity of the theory is a property of the CS theory on a Z-symmetric QCD model, and it is a property of non-commutative theories on a Z-

symmetric QCD model. These properties show that the non-commutativity of the theory is not a product of the standard Coulomb theory in the background of a non-commutative QCD theory.

This paper is organized as follows. In Section 2 we derive the original Z-symmetric QCD model for a coupled field theory. In Section 3 we discuss the non-commutativity of the theory. In Section 4, we show that the non-commutativity of the theory can be realized by the presence of a phase of the matter fields and the non-commutativity of the theory. In Section 5, we show that the non-commutativity of the theory can be realized by the presence of a phase of the matter fields and the non-commutativity of the theory. In Section 6, we show that the non-commutativity of the theory can be realized by the presence of a phase of the matter fields and the non-commutativity of the theory. Finally, we show that the non-commutativity of the model can be realized by the presence of an anomalous behavior of the energy of the fields in the theory.

We have shown that the non-commutativity of the theory is a property of the CS theory on a Z-symmetric QCD model. This property shows that the non-commutativity of the theory can be realized by the presence of a phase of the matter fields and the non-commutativity of the theory. In this paper we have shown that the non-commutativity of the theory on a QCD model is a property of the non-commutative gauge theory. We have shown that the non-commutativity of the theory can be realized by the presence of a phase of the matter fields and the non-commutativity of the theory. In this paper we have shown that the non-commutativity of the theory on

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