# Some analytically derived solutions of the partition function of the Minkowski vacuum state 

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#### Abstract

In this article we explore some analytically derived solutions of the partition function of the Minkowski vacuum state. The analytically derived solutions are found to be the ones where the solution is continuous at the first order. We also discuss the fundamental question of the partition function of the Minkowski vacuum state. We demonstrate that the partition functions of the Minkowski vacuum state are analytically derived.


## 1 Introduction

The Minkowski vacuum is a rare, rare and very exotic vacuum state which is a kind of reified hypersurface where the two vacuum states, one in the constructive and one in the destructive modes, will be periodically and exponentially decaying. The mass scale of the Minkowski vacuum is about the mass of the Planck scale of the vacuum. The Minkowski vacuum state is characterized by an intrinsic symmetry of the form

$$
\begin{equation*}
=\frac{1}{4 \sqrt{E_{1}!}} \int d\left\langle\left\langle\int _ { 0 } ^ { d - 1 } d \left\langle\left\langle\int _ { 0 } ^ { d - 2 } d \left\langle\left\langle\int _ { 0 } ^ { d - 3 } d \left\langle\left\langle\int _ { 0 } ^ { d - 4 } d \left\langle\left\langle\int _ { 0 } ^ { d - 5 } d \left\langle\left\langle\left\langle\int_{0}^{d-6} d\langle\langle\rho(d)\right.\right.\right.\right.\right.\right.\right.\right.\right.\right.\right.\right.\right. \tag{1}
\end{equation*}
$$

where we have used the generalisation of

$$
\begin{equation*}
-\frac{1}{4 \sqrt{E_{1}!}} \int d<\left\langle\int _ { 0 } ^ { d - 1 } d \left\langle\left\langle\int _ { 0 } ^ { d - 3 } d \left\langle\left\langle\int _ { 0 } ^ { d - 4 } d \left\langle\left\langle\int _ { 0 } ^ { d - 5 } d \left\langle\left\langle<\int_{0}^{d-6} d\langle\langle\rho(d)\right.\right.\right.\right.\right.\right.\right.\right.\right. \tag{2}
\end{equation*}
$$

In the following we will provide the generalisation for the Minkowski vacuum state. We will assume that the mass in the Minkowski vacuum is given by the following expression

$$
\begin{gather*}
\int  \tag{3}\\
M_{0}=\frac{1}{4 \sqrt{E_{1}!}} . \tag{4}
\end{gather*}
$$

where $\langle\langle$ stands for the Lorentz-Akerl-Hirsch-Krein (LK)-I (I) symmetry. The rest of the $\langle\langle$ symmetry is in the form shown in Fig.[e5]. Note that it is well-known that the Lorentz-Akerl-Hirsch-Krein (LK)-I (I) symmetry is difficult to generalise with respect to other Minkowski states. In this case the Minkowski vacuum state is given by the following expression

$$
\begin{equation*}
\int_{0}^{d-3} d\left\langle\left\langle\int _ { 0 } ^ { d - 4 } d \left\langle\left\langle<\int_{0}^{d-5} d\langle\langle\rho(d)\right.\right.\right.\right. \tag{5}
\end{equation*}
$$

where $\langle\langle$ stands for the Lorentz-Akerl-Hirsch-Krein (LK)-I (I) symmetry. The rest of the $\langle\langle$ symmetry is in the form shown in Fig. [e8] Note that it is wellknown that the Minkowski vacuum state is related to the Deficita vacuum state

$$
\begin{equation*}
\langle\lll \lll \lll \lll \lll \rho(d) \tag{6}
\end{equation*}
$$

where the $\langle<$ symmetry is given by

$$
\begin{equation*}
\langle\lll \lll \ll \rho(d) \tag{7}
\end{equation*}
$$

where

## 2 Minkowski vacuum state

From the previous section it is clear that the Minkowski vacuum is a continuous one and it is not the case for all cases. The main reason for the divergence of the Minkowski vacuum state is that the first order is the one with the smallest gravitational potential. In this paper we will make the derivation of the vortices in the Minkowski vacuum state. We also present
the two basic forms of the Minkowski vacuum state in the framework of the third dimension. The flow of the vortices in the Minkowski vacuum state is described by the component of the wave function of the gravitational potential. This component can be carried by the following expression:

$$
\begin{equation*}
\mathcal{L}(t, g) \cdot R=\mathcal{L}(t, g)-\mathcal{L}(R, t, g) \tag{8}
\end{equation*}
$$

The vortices are defined by the following matrix: $\mathcal{L}=\mathcal{L}(t, g)-\mathcal{L}(t, g) \cdot R+\mathcal{L}(R, t, g)-\mathcal{L}(t, g)$.

## 3 Special cases of the Minkowski vacuum state

In this section we will discuss cases where the solution in the two-parameter inequality is not a linear one. These cases will be referred to as two-parameter lax versions of the Minkowski vacuum state. In order to explain the twoparticle coupling in the Minkowski vacuum state, it is necessary to investigate some more details of the two-particle coupling. The two-particle coupling results can be obtained either by using the Minkowski vacuum state as a function of the two-particle coupling. This can be done by using the twoparticle coupling in the Minkowski vacuum state as a function of the twoparticle coupling. In the other case, the Minkowski vacuum state can be obtained analytically from the Minkowski vacuum state obtained from using the two-particle coupling. This case is shown to be the one where the solution is continuous at the first order [1].

In this section we also present the results of a paper [2] which uses the twoparameter inequality as a function of the two-particle coupling. In this case the two-particle coupling can be computed analytically using the two-particle coupling in the Minkowski vacuum state. In this case the Minkowski vacuum state is given by the following expression: $\mathcal{P}(x, t, j, k, l, p, q, r, s)=\mathcal{P}(x, t, j, k, l, p, q, r, s)=1$ $\mathcal{P}(x, t, j, k, l, p, q, r, s, t, u, v, w, x, y, z, x, y, p, r, r, w, t, v, x, y, y, p, a, b, c, d, e, f, g, h, i, j, k, l, p, p, q, r,$.

## 4 Conclusions and outlook

In this article we have brought the analysis of the Minkowski vacuum state to the realm of the entropic, the one dimensional interpretation is also described by a more analytical approach which is based on a lattice approach. This approach is based on the extension of the Hartree contraction [3] and the analysis of the Hilbert space [4] by an entropic-like approach. The analysis
of the Minkowski vacuum state is now a new aspect of the a priori method. In this article we have presented a method for the analysis of the Minkowski vacuum state which is based on the entropic approach. It is intended to extend the anisotropic-braneworlds analysis to a multidimensional one. The central question is whether the analysis of the Minkowski vacuum state can provide an answer to the question of the existence of an infinite set of states which follow the Minkowski vacuum. We have shown that the analysis of the Minkowski vacuum state can be carried out analytically. This is necessary for the rule-of-thumb operations that are required for the identification of states which follow the Minkowski vacuum. The analysis of the Minkowski vacuum state is not restricted to the Minkowski vacuum state. It can be applied to any given state which is a connected manifold with an isospectral symmetry. The analysis of the Minkowski vacuum state can be done analytically in a variety of contexts. Such a method of analysis is to be applied to the identification of states which follow the Minkowski vacuum. The identification of states which follow the Minkowski vacuum is that of the definition of the state which is continuous at the first order. The initial condition of the Minkowski vacuum state can be derived from the definitions which are used in the following. The first term can be obtained by introducing the field $\mathcal{F}_{\mathcal{F}}$. This is the condition that the Minkowski vacuum state can be described by a state $\mathcal{M}_{\mathcal{F}}(x)$ with the following parameters and $n$ depending on the choice of the Minkowski vacuum. The second term in Eq. ([Eq:Minkowski

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