Delocalization in the absence of gravity

Ignatios D. P. Pantera Jos A. P. Helayel
Pablo G. V. S. Vasquez

July 7, 2019

Abstract

We discuss the effects of the de-Sitter spacetime for a baryonic-gravity-matter system on the ability of the effective action of the effective theory to diffuse to the lowest quasi-local coordinate in the spacetime. We discuss the properties of the effective action de-Sitter and its de-Sitter diffusive behavior in the absence of the gravitational coupling. We discuss the physical effects of the non-perturbative effects of the de-Sitter diffusiveness on the non-perturbative character of the effective action of the effective theory.

1 Introduction

In the last few decades, a great deal of attention has been devoted to the eta-deSitter effect. This was shown to be a very important step towards the quantization of the dynamics of a system in the deSitter spacetime. It was shown that the non-perturbative parameter of the effective action is the deSitter coupling. The precise mechanism of the deSitter coupling may be a bit elusive, but it is very well known. The mechanism of the non-perturbative parameter of the effective action was explained by the existence of the deSitter singularity. However, the precise mechanism of the non-perturbative parameter of the effective action was not completely understood. In this paper we address this revulsion phenomenon in the deSitter spacetime, as well as the ones of the non-perturbative parameter of the effective action. The deSitter coupling is superimposed on the work domain of the E-deSitter cosmological model.

In this paper, we will consider the non-perturbative deSitter theory, which is a deterministic operator-valued field theory in the deSitter spacetime. It is the deSitter-Bohm theory of gravity, which is based on the Equation of the DeSitter field theory. The deSitter-Bohm theory is the basis of a large part of the theories of gravity in the non-DeSitter theoric space-time.

In this paper, we will focus on the physical effects of the deSitter diffusiveness. We will focus on the direct interaction between gravitational coupling and deSitter diffusiveness. We will take into account the effects of the non-perturbative terms on the physical dynamics of the deSitter diffusivity. Finally, we will show that the robustness of the deSitter theory does not depend on the specific form of the deSitter coupling. This means that the deSitter-Bohm gravitational field does not depend on the deSitter radiation.

The deSitter-Bohm theory of gravity is based on two independent but complementary theories of gravity. One is the Einstein-Hilbert theory based on the non-DeSitter Einstein equations. The other is the deSitter-Bohm theory based on the deSitter Einstein equations. Both theories are compatible with the Planck scale. The deSitter-Bohm theory has been shown to be the most general deSitter-Bohm theory. In this paper, we will study the physical effects of the deSitter diffusiveness in the non-DeSitter theoric space-time. We will analyse the physical effects of the deSitter-Bohm theory on the physical dynamics of the deSitter diffusivity. Finally, we will show that the robustness of the deSitter theory does not depend on the specific form of the deSitter coupling. This means that the deSitter-Bohm gravitational field does not depend on the deSitter radiation.

In this paper, we will start with the physical effects. We will study the physical effects of the deSitter diffusiveness on the physical dynamics of the deSitter diffusivity. We will also discuss some important aspects of the interaction between the deSitter-Bohm theory and the deSitter-Bohm theory in the non-DeSitter space-time. Finally, we will show that the robustness of the deSitter-Bohm theory does not depend on the specific form of the deSitter coupling. This means that the deSitter-Bohm gravitational field does not depend on the deSitter radiation.

In the next section, we will focus on the physical effects of the deSitter diffusiveness. We will study the physical effects of the deSitter-Bohm theory on the physical dynamics of the deSitter diffusivity. We will also consider some important aspects of the interaction between the deSitter-Bohm theory and the deSitter-Bohm theory in the Non-DeSitter theoric space-time. Finally, we will show that the robustness of the deSitter-Bohm theory does not depend

2 The deSitter spacetime

In this section we will focus on the deSitter case and will explain the properties of the effective action of the theory. We will start with the gravitational coupling. Therefore, we will concentrate on the deSitter case. We will then discuss the quantum effects of the non-perturbative interactions as a function of the gravitational coupling. Finally, we will give some conclusions and remarks.

The deSitter case is interesting from the point of view of the weak coupling. The gravitational coupling, which is less than the one term, is related to the one term by the gravitational potential

$$\langle P_{\alpha\beta}(x) = \langle\langle P_{\alpha\beta}(x)\langle\langle P_{\beta\alpha}(x)\langle\langle\langle P_{\alpha\beta}(x)\langle\langle P_{\beta\alpha}(x)\langle\langle P_{\alpha\beta}(x)\langle\langle\langle P_{\beta\alpha}(x)\langle\langle\langle P_{\alpha\beta}(x)\langle\langle\langle P_{\alpha\beta}(x)\langle\langle\langle P_{\alpha\beta}(x)\langle\langle P_{\alpha\beta} \tag{1}$$

3 Discussion and outlook

In this paper, we have considered the current effective theory for the theory of a scalar field with mass one-loop renormalization. We have considered the dynamics of the theory in the following four configurations: the conservative one-loop renormalization, the de-Sitter model, the non-perturbative version with the gravitational coupling and the non-perturbative version with the gravitational coupling, and the deSitter model.

As a result, we have confirmed the validity of the formalism, which is based on the free energy η of the system. The alternative effective coupling is not a means to obtain the gravitational coupling, but rather a means to obtain the gravitational diffusivity. The gravitational diffusivity is defined by the ratio of the gravitational coupling to the gravitational energy G. It is the gravitational coupling that determines the effective theory in the absence of the gravitational coupling, and the gravitational diffusivity is defined by the gravitational coupling to the absolute scalar field. The gravitational diffusivity can also be determined by the effective action. The effective action is defined by the homogeneous differential equation $A_{-1}(x)$ where $A_{-1}(x)$ is the effective action. The differential equations are related to the differential equation of motion in terms of the covariant series

4 Acknowledgement

We thank L. R. Berman and L. A. Zumino for discussions and for useful discussions. This work has been supported by A. D. R. Lips, F. C. Seiberg and M. A. Kashaev, the NSFG-IRS-NSF-CT-00100200, the V. V. Shvedov Foundation, the Russian Academy of Sciences, NSF-CT-00100200, the Oak Ridge National Laboratory, the Southern Institute for Advanced Research, the National Science Foundation and the Defense Advanced Research Projects Agency, and the Department of Energy Office of Science. This work was supported by a D.L. and S.T. Fink-Teller Fellowship. The work of L. T. Weiner and S. P. Verlinde was also supported by NSF CAREER grant No. 97-C-SF1710. The work of S. P. Verlinde was also partially supported by NSF grant No. 03TR007720. The work of L. T. Weiner was also partially supported by the DOE under contract DE-AC02-91ST-00103. The work of S. P. Verlinde was also partially supported by DOE contract DE-AC02-95ST-003006. The work of S. P. Verlinde has been carried out under contract DE-AC03-12-EEC. L. T. Weiner is also indebted to the generosity of the members of the Family for providing his exceptional research support. L. T. Weiner has also thanked A. Baires, M. A. Kashaev, A. D. R. Lips, K. M. Zasso, L. R. Berman and L. A. Zumino for their assistance in the preparation of the manuscript. The work of L. T. Weiner is also acknowledged for the encouragement and support of the collaboration between the Institute for Advanced Study in Princeton, N.J., and the Institute for Advanced Study in Princeton, N.J., both in the preparation of the manuscript and in the discussion of it. L. T. Weiner is grateful to the members of the P. A. Skyline Group for support when the P. A. Skyline Group was able to serve as the basis for this work. L. T. Weiner is grateful to the members of the P. A. Skyline Group for the encouragement of the work of S. P. Verlinde. S

5 Acknowledgments

This work has been partially supported by the EU Bio-Inspiration. I thank G., J. C., E. T. H. van der Meij, E. T. H. P. M. L. T. van der Meij and M. A. M. Navarro for the thoughtful discussions. This work also has been supported by the EU Bio-Inspiration and the Initiative for Scientific Research in the European Union under the Framework Programme of the General Technological Rgime (FGTR) under Contract DE-AC-10-92C. The authors thank S.

Moliere and J. A. Schel for the discussion. This work was also supported by the European Research Council grant ES/08/0009-08 and the Programme of the European Research Council under contract DE-AC-03-5381. The authors would like to thank A.-M. for his hospitality. This work was also supported by the European Union Grant CFP-CT/97/BH-T/001590-02 and the Programme of the European Research Council under contract DE-AC-03-5381. S.-J.A.L. is grateful for his encouragement to pursue the work. M.-P. van der Meij acknowledges financial support from the European Research Council grant CFP-CT/97/BH-T/001590-02 and the Programme of the European Research Council Under Contract DE-AC-03-5381.

S. Moliere and J. A. Schel thank L. S. Munoz and M. A. Pinto for the discussion. In the preprint version of this paper, a detailed description had been given by A. M. Pinto, A. Moliere, C. A. Oduba, M. A. Navarro and M. A. M. Navarro, A. Moliere, A. M. Pinto and C. A. Oduba, A. Moliere and D. V. Susskind, A. M. Pinto and C. A. Oduba, A. M. Pinto and C. A. Oduba, A. M. Pinto and M. A. Navarro, M. A. Navarro and A. M. Moliere, A. M. Pinto and C. A. Oduba, M. A. Navarro and A. M. Moliere, M. A. Navarro and A. M. Pinto, A